



Division of Strength of Materials and Structures

Faculty of Power and Aeronautical Engineering

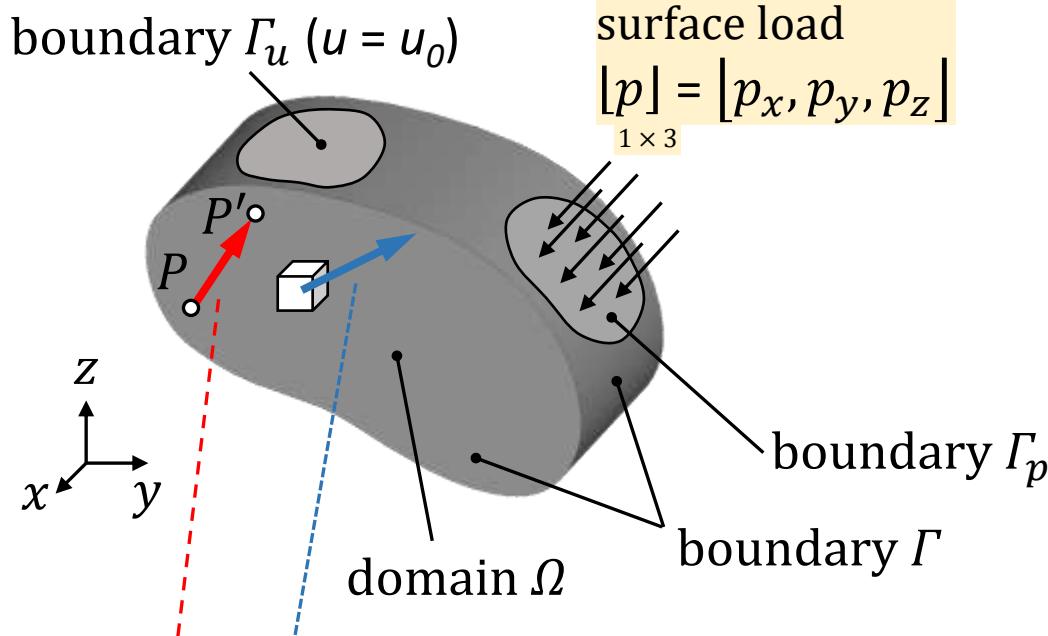


# Finite element method (FEM1)

Lecture 2A. The boundary value problem of solid mechanics  
in the FEM approach

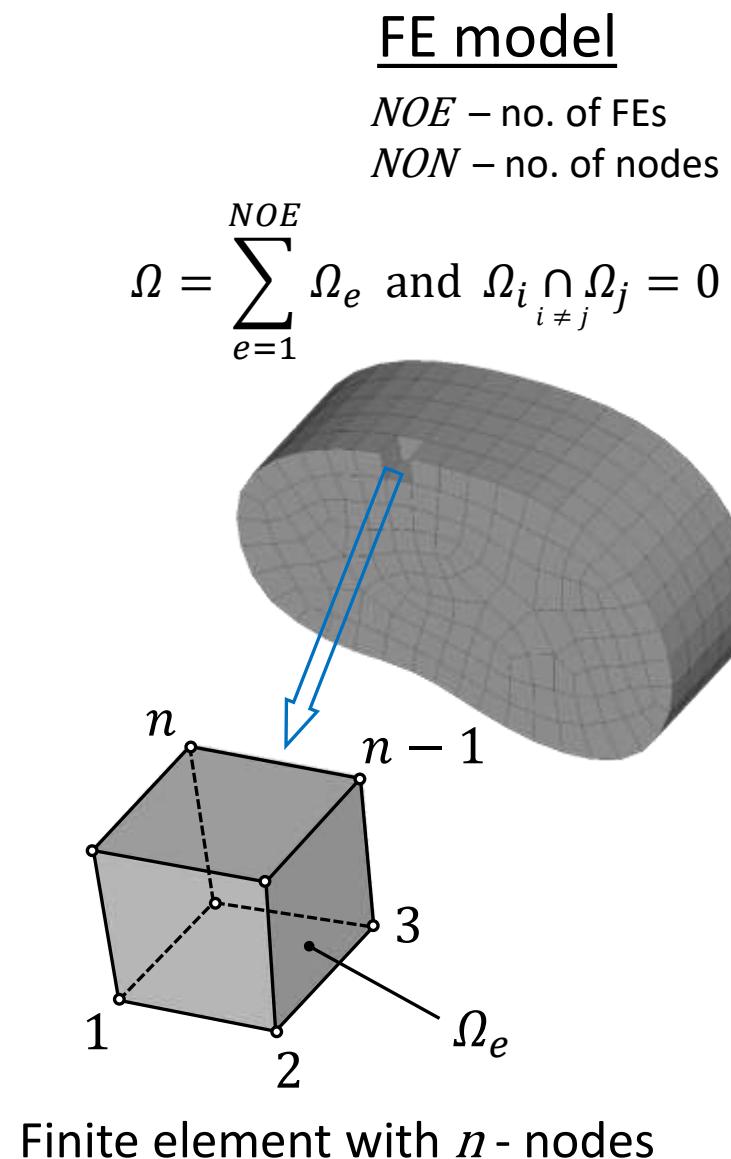
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# Boundary value problem of solid body mechanics

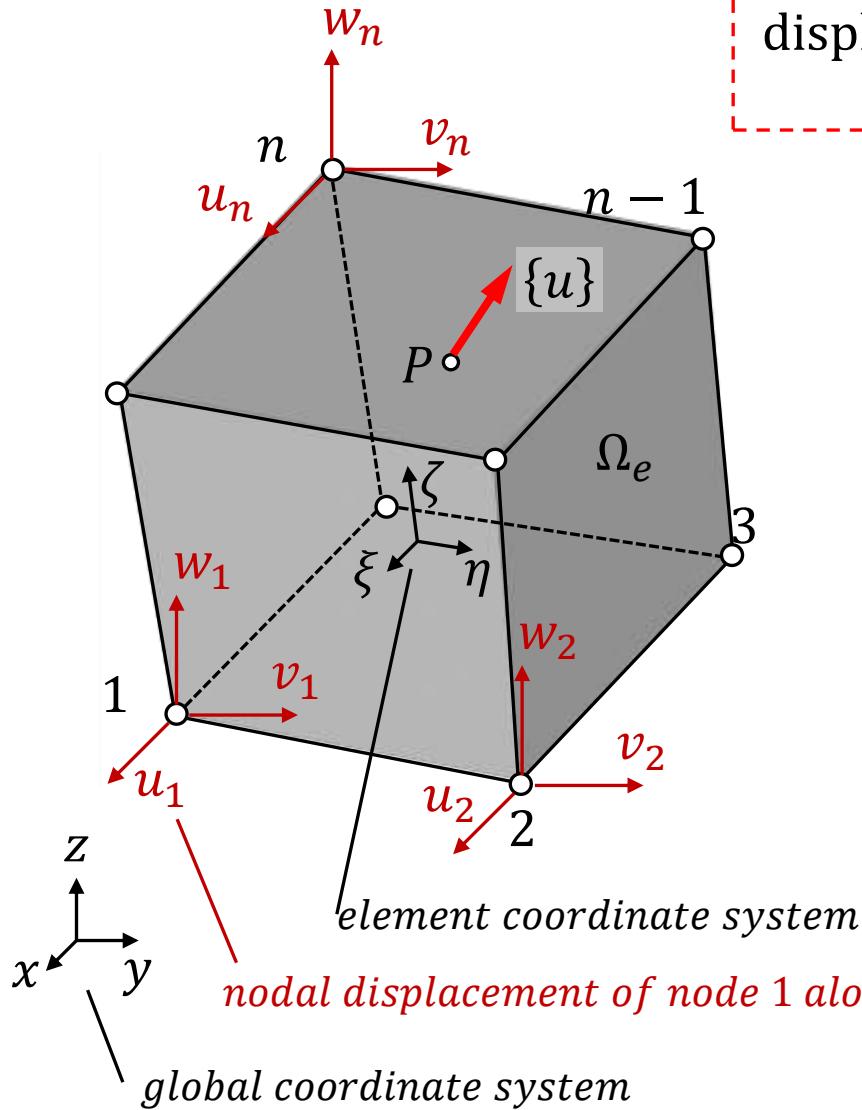


**UNKNOWN FUNCTION**

Displacement vector  $\{u\} = \begin{Bmatrix} u(x, y, z) \\ v(x, y, z) \\ w(x, y, z) \end{Bmatrix}$



# Nodal approximation inside the finite element with n - nodes



displacement vector  $\{u\} = [N(\xi, \eta, \zeta)]\{q\}_e$

$3 \times 1$	$3 \times n_e$	$n_e \times 1$
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$[N(\xi, \eta, \zeta)]$  – matrix of shape functions  
 $3 \times n_e$

$$n_e = n \cdot n_p$$

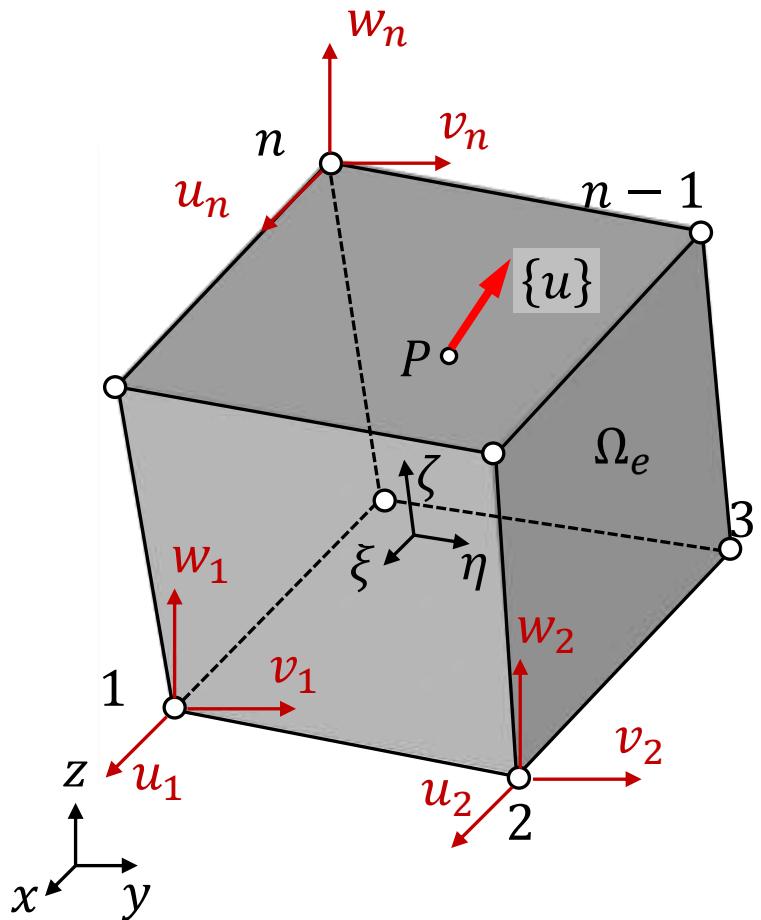
$n_e$  – no. of degrees of freedom in FE

$n_p$  – no. of degrees of freedom per node

$$\{q\}_e = \left\{ \begin{array}{c} u_1 \\ v_1 \\ w_1 \\ \vdots \\ u_n \\ v_n \\ w_n \end{array} \right\}_{n_e \times 1}$$

– local vector of nodal parameters

# Matrix of shape functions



Shape function matrix:

$$[N(\xi, \eta, \zeta)] = \begin{bmatrix} N_1 & 0 & 0 & N_2 & 0 & 0 & N_n & 0 & 0 \\ 0 & N_1 & 0 & 0 & N_2 & 0 & \dots & 0 & N_n \\ 0 & 0 & N_1 & 0 & 0 & N_2 & 0 & 0 & N_n \end{bmatrix}_{3 \times n_e}$$

Nodal approximation:

$$\{u\} = [N(\xi, \eta, \zeta)]\{q\}_e$$

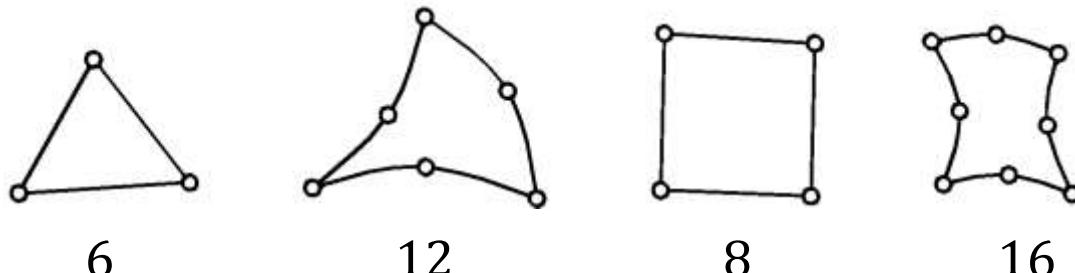
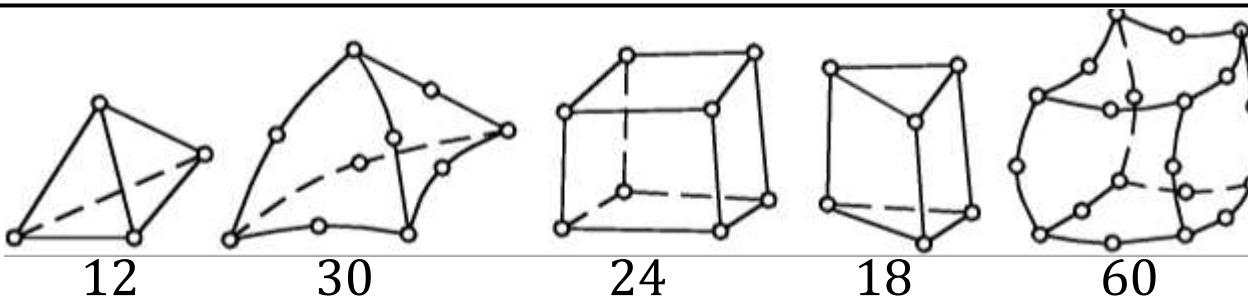
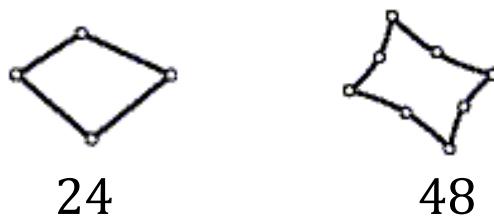
$3 \times 1$        $3 \times n_e$        $n_e \times 1$

$$\begin{aligned} u &= N_1 \cdot u_1 + N_2 \cdot u_2 + \dots + N_n \cdot u_n \\ v &= N_1 \cdot v_1 + N_2 \cdot v_2 + \dots + N_n \cdot v_n \\ w &= N_1 \cdot w_1 + N_2 \cdot w_2 + \dots + N_n \cdot w_n \end{aligned}$$

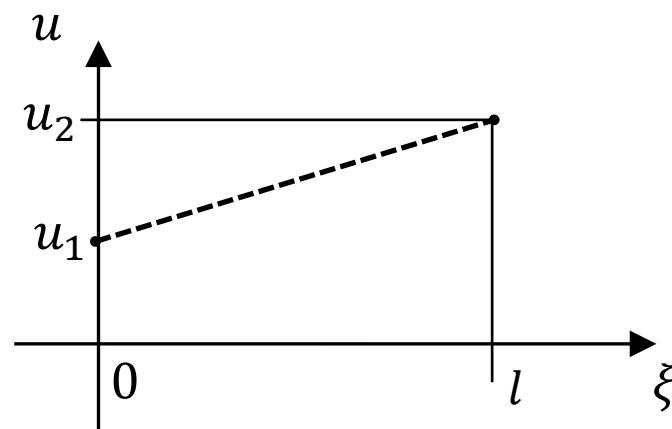
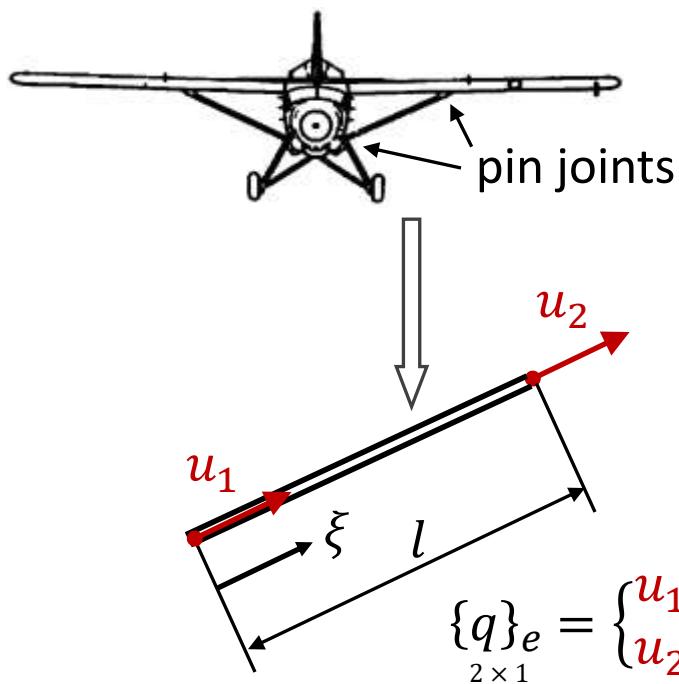
Vector of element nodal parameters:

$$\{q\}_e = \left\{ \begin{array}{c} u_1 \\ v_1 \\ w_1 \\ \vdots \\ u_n \\ v_n \\ w_n \end{array} \right\}_e$$

# Examples of finite elements

Type	$n_e$ – number of degrees of freedom in FE
rods	
2D	
3D	
shell	

# Example 1: shape functions for a finite element representing a strut



linear function:

$$u(\xi) = \frac{u_2 - u_1}{l} \xi + u_1$$

$$\begin{aligned} u(\xi) &= \frac{u_2 - u_1}{l} \xi + u_1 = \frac{u_2}{l} \xi - \frac{u_1}{l} \xi + u_1 = \left(1 - \frac{\xi}{l}\right) u_1 + \frac{\xi}{l} u_2 = \\ &= N_1(\xi) \cdot u_1 + N_2(\xi) \cdot u_2 = [N_1, N_2] \begin{Bmatrix} u_1 \\ u_2 \end{Bmatrix}_e = [N(\xi)] \{q\}_e \end{aligned}$$

shape functions:

$$N_1(\xi) = 1 - \frac{\xi}{l}$$

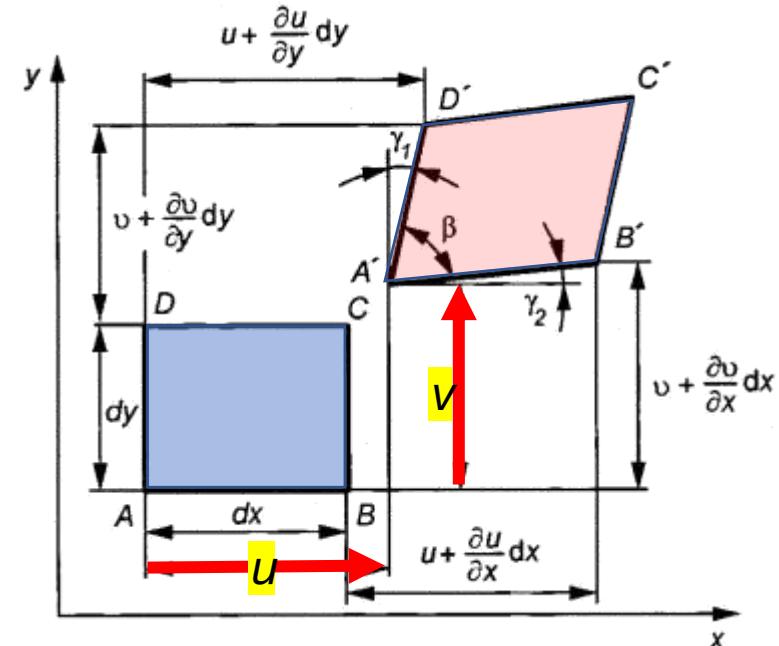
$$N_2(\xi) = \frac{\xi}{l}$$

# Strain components

normal strains:

$$\varepsilon_x = \frac{(A'B')_x - AB}{AB} = \frac{(dx + u + \frac{\partial u}{\partial x} dx - u) - dx}{dx} = \boxed{\frac{\partial u}{\partial x}}$$

$$\varepsilon_y = \boxed{\frac{\partial v}{\partial y}} ; \quad \varepsilon_z = \boxed{\frac{\partial w}{\partial z}}$$



shear strains:

$$\gamma_{xy} = \frac{\pi}{2} - \beta = \gamma_1 + \gamma_2$$

$$\gamma_1 \cong \tan \gamma_1 = \frac{(A'D')_x}{(A'D')_y} = \frac{u + \frac{\partial u}{\partial y} dy - u}{dy + v + \frac{\partial v}{\partial y} dy - v} = \frac{\frac{\partial u}{\partial y}}{1 + \frac{\partial v}{\partial y}} = \frac{\frac{\partial u}{\partial y}}{1 + \varepsilon_y} = \frac{\partial u}{\partial y}$$

$$\gamma_2 \cong \frac{\partial v}{\partial x} \rightarrow \boxed{\gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}}$$

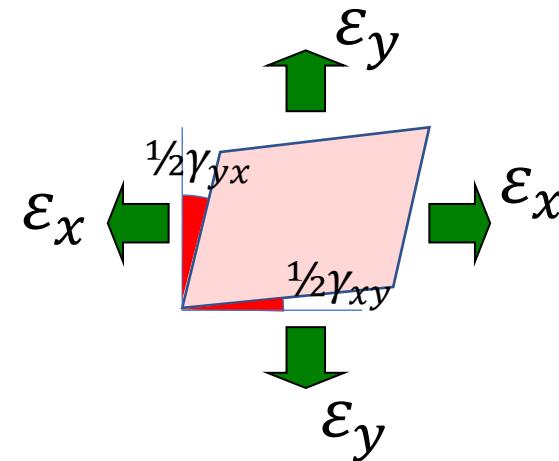
*small deformations:  $\varepsilon_y \ll 1$*

$$\boxed{\gamma_{yz} = \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y}} ; \quad \boxed{\gamma_{zx} = \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z}} ; \quad \gamma_{ij} = \gamma_{ji}$$

# Strain tensor. Vector of strain components

strain tensor:

$$\boldsymbol{\varepsilon} = \begin{bmatrix} \varepsilon_x & \gamma_{xy}/2 & \gamma_{xz}/2 \\ \gamma_{yx}/2 & \varepsilon_y & \gamma_{yz}/2 \\ \gamma_{zx}/2 & \gamma_{zy}/2 & \varepsilon_z \end{bmatrix}_{3 \times 3}$$



vector of strain components:

$$\left\{ \varepsilon \right\}_{6 \times 1} = \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \varepsilon_z \\ \gamma_{xy} \\ \gamma_{yz} \\ \gamma_{zx} \end{Bmatrix} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 & 0 \\ 0 & \frac{\partial}{\partial y} & 0 \\ 0 & 0 & \frac{\partial}{\partial z} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial z} & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} & 0 & \frac{\partial}{\partial x} \end{bmatrix}_{6 \times 3} \begin{Bmatrix} u \\ v \\ w \end{Bmatrix}_{3 \times 1} = [R] \{u\}_{6 \times 3} \quad ; \quad [\varepsilon]_{1 \times 6} = [u]_{1 \times 3} [R]_{3 \times 6}^T$$

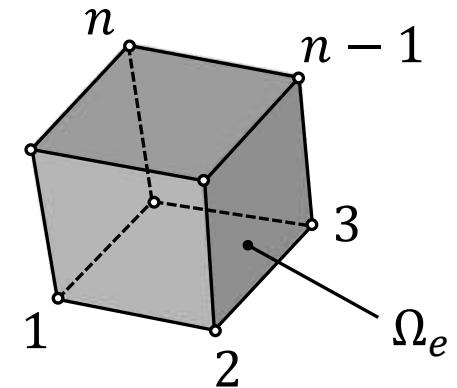
*gradient matrix*

# Strain – displacement matrix of a finite element

nodal approximation in a finite element:

$$\{u\} = [N(\xi, \eta, \zeta)]\{q\}_e$$

$3 \times 1 \quad 3 \times n_e \quad n_e \times 1$



vector of strain components in a finite element:

$$\{\varepsilon\} = [R]\{u\} = [R][N]\{q\}_e = [B]\{q\}_e$$

$6 \times 1 \quad 6 \times 3 \quad 3 \times 1 \quad 6 \times 3 \quad 3 \times n_e \quad n_e \times 1 \quad 6 \times n_e \quad n_e \times 1$

$$[\varepsilon] = [q]_e [B]^T$$

$1 \times 6 \quad 1 \times n_e \quad n_e \times 6$

$$[B] = [R][N] \quad - \text{strain-displacement matrix}$$

$6 \times n_e \quad 6 \times 3 \quad 3 \times n_e$

# Stress components

normal stresses:

$$\sigma_x \ ; \ \sigma_y \ ; \ \sigma_z$$

positive value - tension, negative value - compression

shear stress components:

$$\tau_{xy} \ ; \ \tau_{yz} \ ; \ \tau_{zx} \ ; \ \tau_{ij} = \tau_{ji}$$

equivalent stresses:

Von Mises stress:

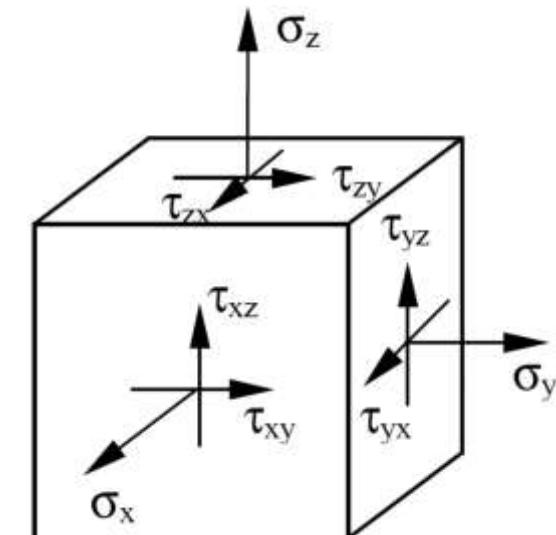
$$\sigma_{EQV} = \sqrt{\frac{1}{2} \left( (\sigma_x - \sigma_y)^2 + (\sigma_y - \sigma_z)^2 + (\sigma_z - \sigma_x)^2 \right) + 3(\tau_{xy}^2 + \tau_{yz}^2 + \tau_{zx}^2)}$$

Tresca stress:

$$\sigma_{INT} = \sigma_1 - \sigma_3 = 2\tau_{max}$$

the first  
principal stress

the third  
principal stress



maximum shear stress

# Stress tensor. Vector of stress components

stress tensor:

$$\sigma = \begin{bmatrix} \sigma_x & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \sigma_y & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \sigma_z \end{bmatrix}_{3 \times 3} \equiv \begin{bmatrix} \sigma_1 & 0 & 0 \\ 0 & \sigma_2 & 0 \\ 0 & 0 & \sigma_3 \end{bmatrix}$$

*in coordinate system x, y, z*

*in the principal coordinate system*

vector of stress components:

$$\{\sigma\} = \begin{Bmatrix} \sigma_x \\ \sigma_y \\ \sigma_z \\ \tau_{xy} \\ \tau_{yz} \\ \tau_{zx} \end{Bmatrix}_{6 \times 1}$$

# Constitutive matrix

linear isotropic material (Hooke's law):

$$\{\sigma\} = [D] \{\varepsilon\}$$

6 × 1      6 × 6      6 × 1



constitutive matrix:

$$[D]_{6 \times 6} = \frac{E}{(1+\nu)(1-2\nu)} \begin{bmatrix} 1-\nu & \nu & \nu & 0 & 0 & 0 \\ \nu & 1-\nu & \nu & 0 & 0 & 0 \\ \nu & \nu & 1-\nu & 0 & 0 & 0 \\ 0 & 0 & 0 & 0.5-\nu & 0 & 0 \\ 0 & 0 & 0 & 0 & 0.5-\nu & 0 \\ 0 & 0 & 0 & 0 & 0 & 0.5-\nu \end{bmatrix}$$

$E$  – Young's modulus,  
 $\nu$  – Poisson's ratio

## Example 2: uniaxial tensile test

$$\sigma_x = \frac{F}{A_0} ; \quad \varepsilon_x = \frac{L - L_0}{L_0} ; \quad \varepsilon_y = \varepsilon_z = \varepsilon_T$$

elastic strain Energy:  $U = \frac{1}{2} \sigma_x \varepsilon_x A_0 L_0$

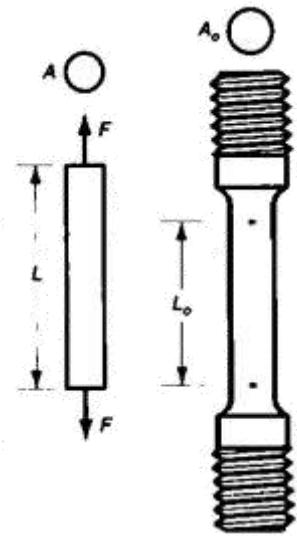
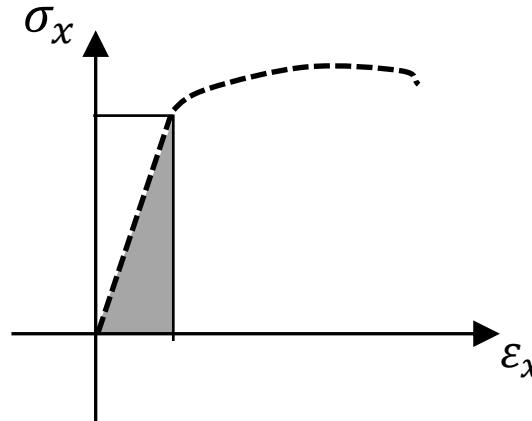
$$\{\sigma\} = [D] \{\varepsilon\}$$

$$\begin{Bmatrix} \sigma_x \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{Bmatrix} = \frac{E}{(1+\nu)(1-2\nu)} \begin{bmatrix} 1-\nu & \nu & \nu & 0 & 0 & 0 \\ \nu & 1-\nu & \nu & 0 & 0 & 0 \\ \nu & \nu & 1-\nu & 0 & 0 & 0 \\ 0 & 0 & 0 & 0.5-\nu & 0 & 0 \\ 0 & 0 & 0 & 0 & 0.5-\nu & 0 \\ 0 & 0 & 0 & 0 & 0 & 0.5-\nu \end{bmatrix} \begin{Bmatrix} \varepsilon_x \\ \varepsilon_T \\ \varepsilon_T \\ 0 \\ 0 \\ 0 \end{Bmatrix}$$

2nd equation:  $0 = \frac{E}{(1+\nu)(1-2\nu)} (\nu \varepsilon_x + (1-\nu) \varepsilon_T + \nu \varepsilon_T) \rightarrow \boxed{\varepsilon_T = -\nu \varepsilon_x}$

1st equation:

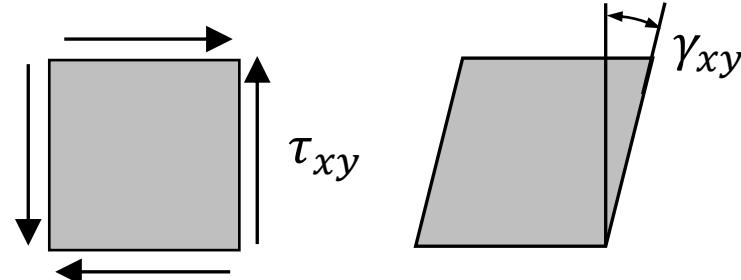
$$\sigma_x = \frac{E}{(1+\nu)(1-2\nu)} ((1-\nu) \varepsilon_x + \nu \varepsilon_T + \nu \varepsilon_T) = \frac{E}{(1-\nu-2\nu^2)} ((1-\nu) \varepsilon_x - \nu^2 \varepsilon_x - \nu^2 \varepsilon_x) \rightarrow \boxed{\sigma_x = E \varepsilon_x}$$



## Example 3: pure shear

$$\tau_{xy} ; \gamma_{xy}$$

$$\begin{matrix} \{\sigma\} = [D] \{\varepsilon\} \\ 6 \times 1 \quad 6 \times 6 \quad 6 \times 1 \end{matrix}$$



$$\begin{Bmatrix} 0 \\ 0 \\ 0 \\ \tau_{xy} \\ 0 \\ 0 \end{Bmatrix} = \frac{E}{(1+\nu)(1-2\nu)} \begin{bmatrix} 1-\nu & \nu & \nu & 0 & 0 & 0 \\ \nu & 1-\nu & \nu & 0 & 0 & 0 \\ \nu & \nu & 1-\nu & 0 & 0 & 0 \\ 0 & 0 & 0 & 0.5-\nu & 0 & 0 \\ 0 & 0 & 0 & 0 & 0.5-\nu & 0 \\ 0 & 0 & 0 & 0 & 0 & 0.5-\nu \end{bmatrix} \begin{Bmatrix} 0 \\ 0 \\ 0 \\ \gamma_{xy} \\ 0 \\ 0 \end{Bmatrix}$$

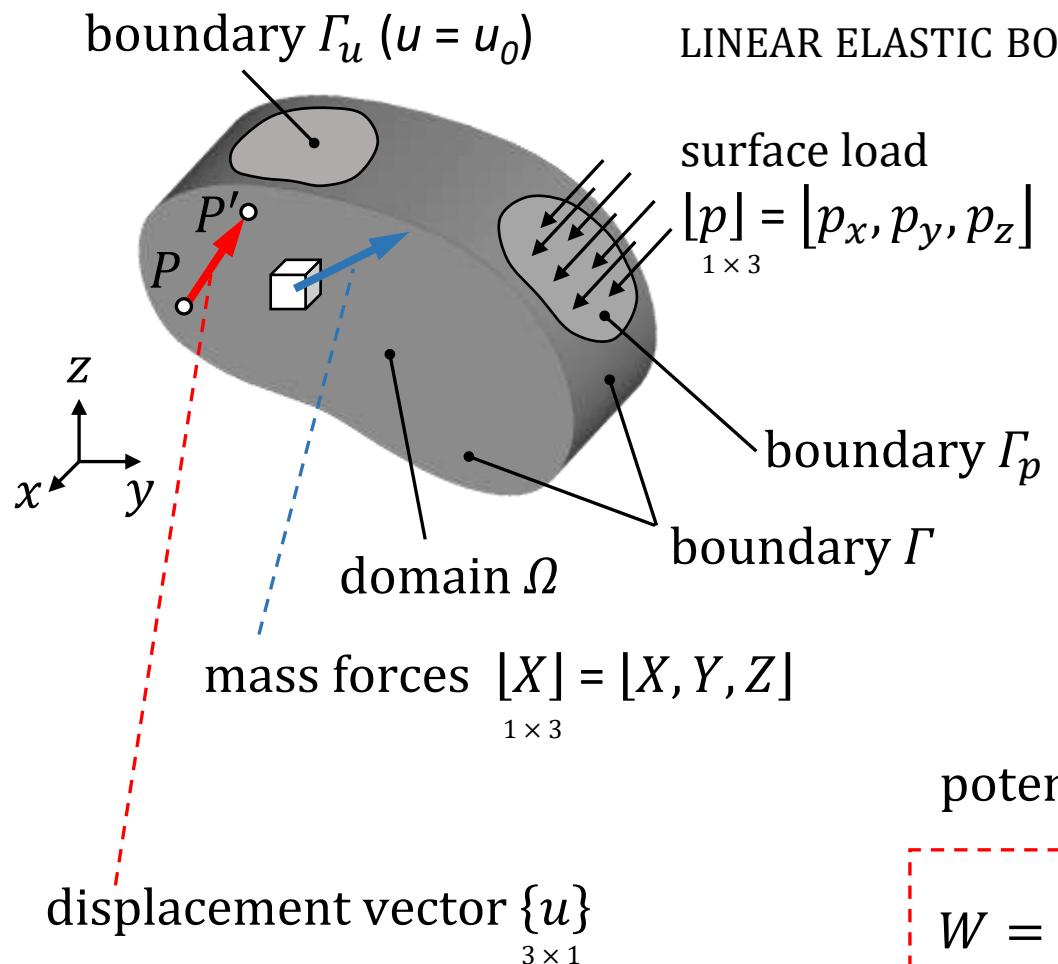
4th equation:

$$\tau_{xy} = \frac{E}{(1+\nu)(1-2\nu)} (0.5 - \nu) \gamma_{xy} = \frac{E}{2(1+\nu)(0.5-\nu)} (0.5 - \nu) \gamma_{xy} = \frac{E}{2(1+\nu)} \gamma_{xy} \rightarrow$$

$\tau_{xy} = G \gamma_{xy}$

$$G = \frac{E}{2(1+\nu)} - \text{Kirchoff's modulus (shear modulus)}$$

# Elastic strain energy. Potential energy of loading



elastic strain energy:

$$U = \frac{1}{2} \int_{\Omega} [\varepsilon]\{\sigma\} d\Omega$$

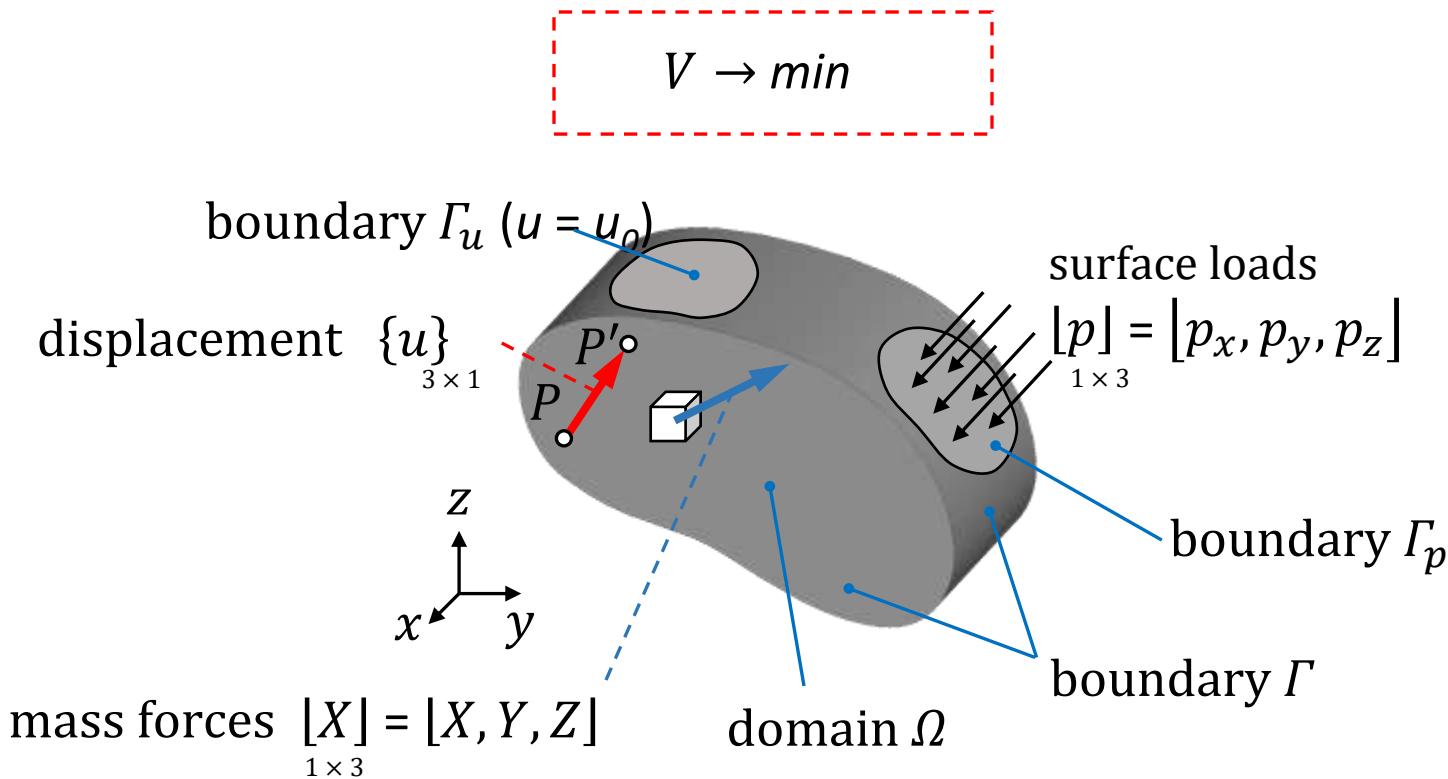
potential energy of loading:

$$W = \int_{\Omega} [X]\{u\} d\Omega + \int_{\Gamma_p} [p]\{u\} d\Gamma_p$$

# Minimum total potential energy principle

$$\text{total potential energy: } V = U - W$$

The displacement field  $\{u\}$  that represents solution of the problem fulfils displacement boundary conditions on  $\Gamma_u$  and minimizes the total potential energy  $V$ .



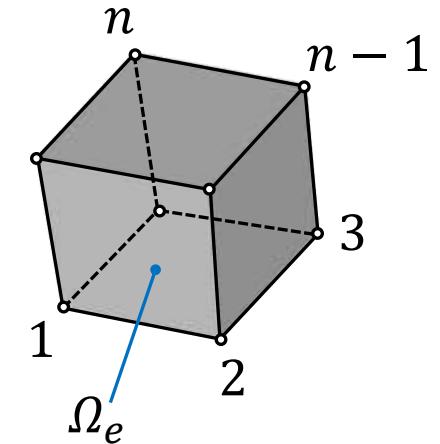
# Elastic strain energy in a finite element. Local stiffness matrix

$\{q\}_e$  - local vector of nodal parameters  
 $n_e \times 1$

elastic strain energy in a finite element:

$$U_e = \frac{1}{2} \int_{\Omega_e} [\varepsilon] \{\sigma\} d\Omega_e = \frac{1}{2} \begin{bmatrix} q \end{bmatrix}_e \int_{\Omega_e} [B]^T [D] [B] d\Omega_e \quad \{q\}_e = \frac{1}{2} \begin{bmatrix} q \end{bmatrix}_e [k]_e \{q\}_e$$

$\begin{matrix} 1 \times 6 & 6 \times 1 \\ \uparrow & \uparrow \\ \{q\}_e & = [B] \{q\}_e \end{matrix}$ 
 $\begin{matrix} 6 \times 1 & 6 \times 6 & 6 \times 1 \\ \uparrow & \uparrow & \uparrow \\ [\varepsilon] & = [q]_e [B]^T & \{q\}_e = [B] \{q\}_e \end{matrix}$

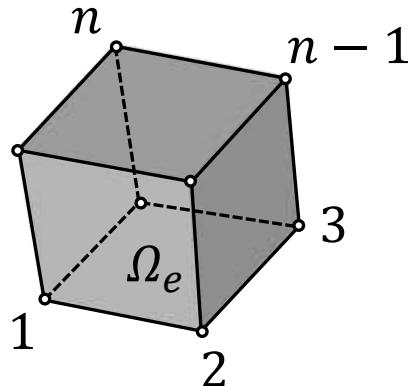


local stiffness matrix:

$$[k]_e = \int_{\Omega_e} [B]^T [D] [B] d\Omega_e$$

# Elastic strain energy in a finite element

local notation:



$n$  – no. of nodes per FE

$n_p$  – no. of nodal parameters per node

no. of degrees of freedom in FE:

$$n_e = n \cdot n_p$$

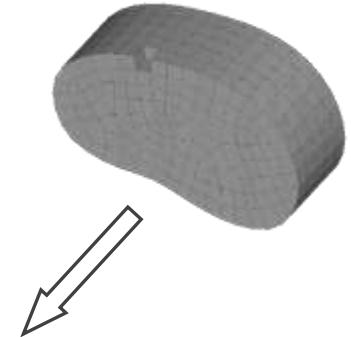
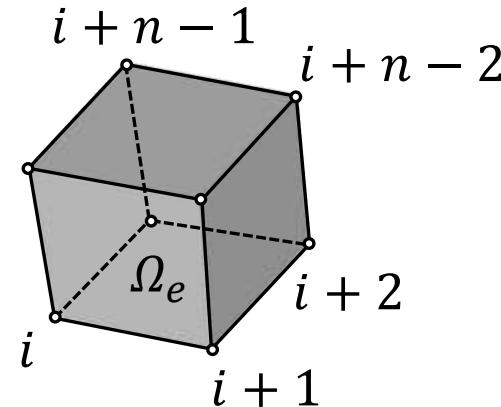
$\{q\}_e$  - local vector of nodal parameters  
 $n_e \times 1$

$$U_e = \frac{1}{2} \underset{1 \times n_e}{[q]_e} \underset{n_e \times n_e}{[k]_e} \underset{n_e \times 1}{\{q\}_e}$$



local stiffness matrix

global notation:



$NON$  – no. of nodes

$n_p$  – no. of nodal parameters per node

no. of degrees of freedom :

$$NDOF = NON \cdot n_p$$

$\{q\}$  - global vector of nodal parameters

$NDOF \times 1$

$$U_e = \frac{1}{2} \underset{1 \times NDOF}{\cdot [q]} \underset{NDOF \times NDOF}{\cdot [k]_e^*} \underset{NDOF \times 1}{\cdot \{q\}}$$



extended local stiffness matrix.

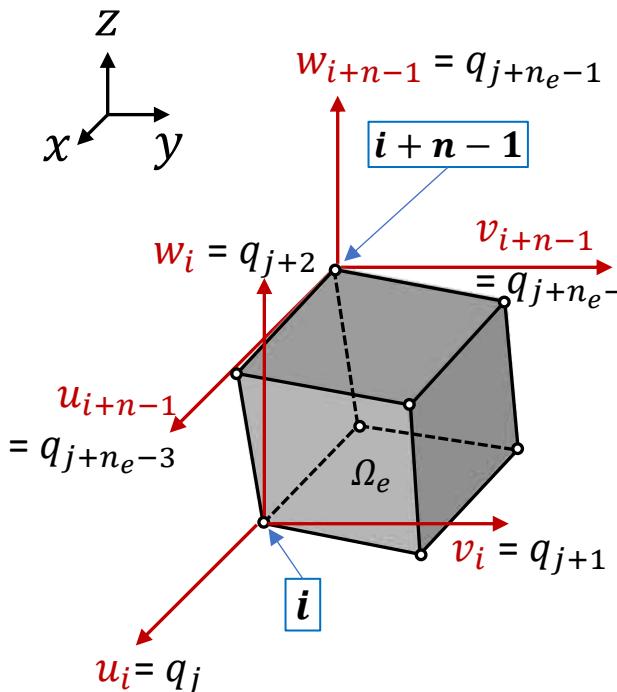
# Extended local stiffness matrix of a finite element

$$\{q\} = \begin{Bmatrix} q_1 \\ q_2 \\ \vdots \\ q_j \\ \vdots \\ q_{NDOF} \end{Bmatrix}_{NDOF \times 1}$$

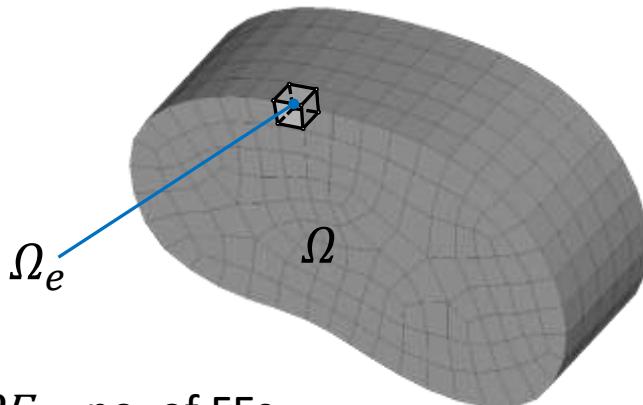
$[k]_e^* =$

	1	2	...	$j-1$	$j$	$j+1$	...	$j+n_e-1$	$j+n_e$	...	$NDOF$
1	0	0	...	0	0	0	...	0	0	...	0
2	0	0	...	0	0	0	...	0	0	...	0
...	...	...	...	0	0	0	...	0	0	...	0
$j-1$	0	0	0	0	0	0	...	0	0	...	0
$j$	0	0	0	0	$k_{11}$	$k_{12}$	...	$k_{1n_e}$	0	...	0
$j+1$	0	0	0	0	$k_{21}$	$k_{22}$	...	$k_{2n_e}$	0	...	0
...	...	...	...	...	...	...	...	...	0	...	0
$j+n_e$	0	0	0	0	$k_{n_e 1}$	$k_{n_e 2}$	...	$k_{n_e n_e}$	0	...	0
...	...	...	...	...	0	0	0	0	0	...	0
$NDOF$	0	0	0	0	0	0	0	0	0	0	0

(assumed ascending order of components)



# Elastic strain energy in a FE model. Global stiffness matrix



$NOE$  – no. of FEs

$NDOF$  – no. of degrees of freedom

$$\Omega = \sum_{e=1}^{NOE} \Omega_e \rightarrow$$

$$U = \sum_{e=1}^{NOE} U_e$$

$\{q\}$  - global vector of nodal parameters  
 $NDOF \times 1$

elastic strain energy in a finite element model:

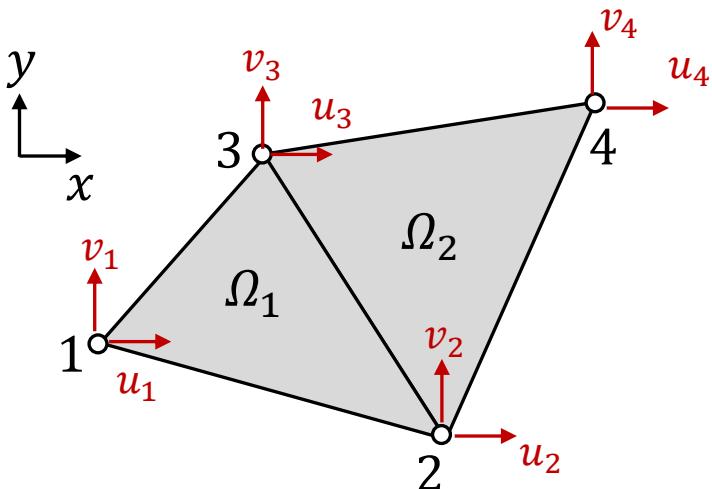
$$U = \sum_{e=1}^{NOE} U_e = \sum_{e=1}^{NOE} \frac{1}{2} \cdot [q] \cdot [k]_e^* \cdot \{q\}$$

$$= \frac{1}{2} \cdot [q] \cdot [K] \cdot \{q\}$$

↑  
global stiffness matrix:

$$[K] = \sum_{e=1}^{NOE} [k]_e^*$$

## Example 4: global stiffness matrix of a 2D model with two 3-node triangles



global notation:

$$\begin{aligned} NOE &= 2 \\ NON &= 4 \\ n &= 3 \\ n_p &= 2 \quad ; \quad (u, v) \\ n_e &= n \cdot n_p = 6 \\ NDOF &= NON \cdot n_p = 8 \end{aligned}$$

$$\{q\} = \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \\ q_5 \\ q_6 \\ q_7 \\ q_8 \end{Bmatrix}_{8 \times 1} = \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{Bmatrix}$$

local notation:

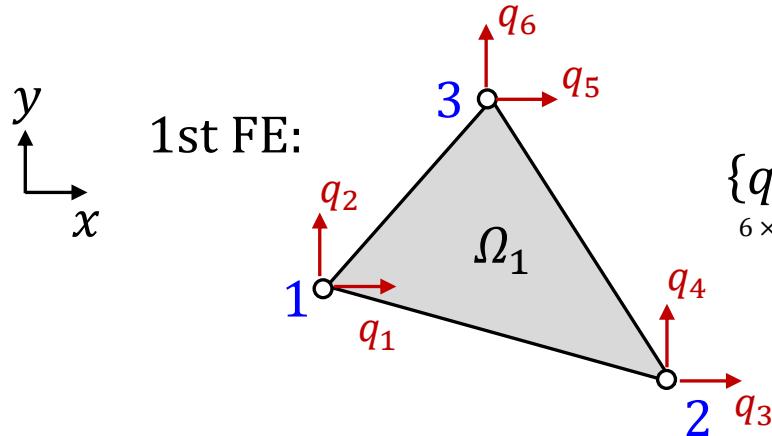
$$\{q\}_1 = \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \\ q_5 \\ q_6 \end{Bmatrix}_{6 \times 1} = \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{Bmatrix}_1$$

A diagram of triangle  $\Omega_1$  with vertices 1, 2, and 3. The local degrees of freedom are labeled as  $q_1, q_2, q_3, q_4, q_5, q_6$ . Vertex 1 is at the bottom-left, vertex 2 at the bottom-right, and vertex 3 at the top. Arrows indicate the direction of each degree of freedom.

$$\{q\}_2 = \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \\ q_5 \\ q_6 \end{Bmatrix}_{6 \times 1} = \begin{Bmatrix} u_2 \\ v_2 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{Bmatrix}_2$$

A diagram of triangle  $\Omega_2$  with vertices 2, 3, and 4. The local degrees of freedom are labeled as  $q_1, q_2, q_3, q_4, q_5, q_6$ . Vertex 1 is at the bottom, vertex 2 is at the top-left, and vertex 3 is at the top-right. Arrows indicate the direction of each degree of freedom.

## Example 4: global stiffness matrix of a 2D model with two 3-node triangles



$$\{q\}_1 = \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \\ q_5 \\ q_6 \end{Bmatrix}_{6 \times 1} = \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{Bmatrix}_1$$

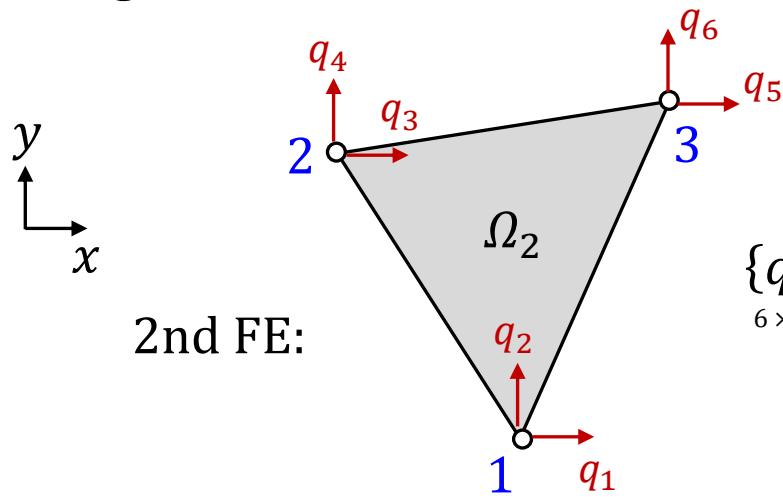
$\{q\} = \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{Bmatrix}_{8 \times 1}$

	1	2	3	4	5	6	
1	$a_1$	$b_1$	$c_1$	$d_1$	$e_1$	$f_1$	
2	$b_1$	$g_1$	$h_1$	$i_1$	$j_1$	$k_1$	
3	$c_1$	$h_1$	$l_1$	$m_1$	$n_1$	$o_1$	
4	$d_1$	$i_1$	$m_1$	$p_1$	$r_1$	$s_1$	
5	$e_1$	$j_1$	$n_1$	$r_1$	$t_1$	$\bar{u}_1$	
6	$f_1$	$k_1$	$o_1$	$s_1$	$\bar{u}_1$	$\bar{w}_1$	

$$[k]_1^* =$$

1	2	3	4	5	6	7	8	
1	$a_1$	$b_1$	$c_1$	$d_1$	$e_1$	$f_1$	0	0
2	$b_1$	$g_1$	$h_1$	$i_1$	$j_1$	$k_1$	0	0
3	$c_1$	$h_1$	$l_1$	$m_1$	$n_1$	$o_1$	0	0
4	$d_1$	$i_1$	$m_1$	$p_1$	$r_1$	$s_1$	0	0
5	$e_1$	$j_1$	$n_1$	$r_1$	$t_1$	$\bar{u}_1$	0	0
6	$f_1$	$k_1$	$o_1$	$s_1$	$\bar{u}_1$	$\bar{w}_1$	0	0
7	0	0	0	0	0	0	0	0
8	0	0	0	0	0	0	0	0

## Example 4: global stiffness matrix of a 2D model with two 3-node triangles



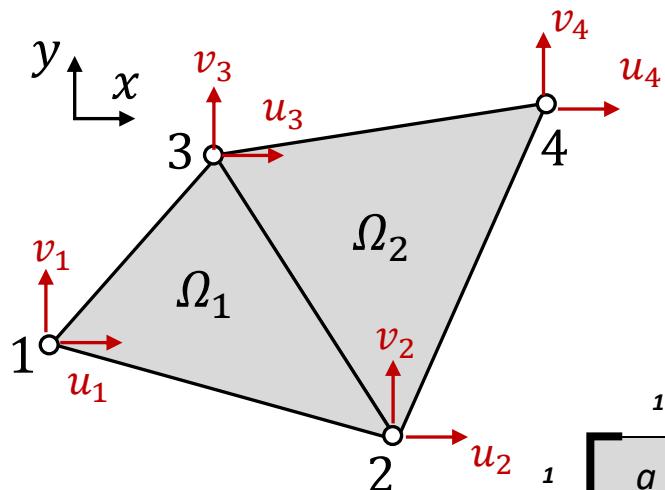
$$\{q\}_2 = \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \\ q_5 \\ q_6 \end{Bmatrix}_{6 \times 1} = \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{Bmatrix}_2 \quad \{q\} = \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{Bmatrix}_{8 \times 1}$$

	1	2	3	4	5	6
1	$a_2$	$b_2$	$c_2$	$d_2$	$e_2$	$f_2$
2	$b_2$	$g_2$	$h_2$	$i_2$	$j_2$	$k_2$
3	$c_2$	$h_2$	$l_2$	$m_2$	$n_2$	$o_2$
4	$d_2$	$i_2$	$m_2$	$p_2$	$r_2$	$s_2$
5	$e_2$	$j_2$	$n_2$	$r_2$	$t_2$	$\bar{u}_2$
6	$f_2$	$k_2$	$o_2$	$s_2$	$\bar{u}_2$	$\bar{w}_2$

$$[k]_2^* =$$

1	2	3	4	5	6	7	8	
1	0	0	0	0	0	0	0	
2	0	0	0	0	0	0	0	
3	0	0	$a_2$	$b_2$	$c_2$	$d_2$	$e_2$	$f_2$
4	0	0	$b_2$	$g_2$	$h_2$	$i_2$	$j_2$	$k_2$
5	0	0	$c_2$	$h_2$	$l_2$	$m_2$	$n_2$	$o_2$
6	0	0	$d_2$	$i_2$	$m_2$	$p_2$	$r_2$	$s_2$
7	0	0	$e_2$	$j_2$	$n_2$	$r_2$	$t_2$	$\bar{u}_2$
8	0	0	$f_2$	$k_2$	$o_2$	$s_2$	$\bar{u}_2$	$\bar{w}_2$

## Example 4: global stiffness matrix of a 2D model with two 3-node triangles



$$[K] = [k]^*_1 + [k]^*_2 =$$

$$\{q\} = \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{Bmatrix}_{8 \times 1}$$

	1	2	3	4	5	6	7	8
1	$a_1$	$b_1$	$c_1$	$d_1$	$e_1$	$f_1$	0	0
2	$b_1$	$g_1$	$h_1$	$i_1$	$j_1$	$k_1$	0	0
3	$c_1$	$h_1$	$l_1 + a_2$	$m_1 + b_2$	$n_1 + c_2$	$o_1 + d_2$	$e_2$	$f_2$
4	$d_1$	$i_1$	$m_1 + b_2$	$p_1 + g_2$	$r_1 + h_2$	$s_1 + i_2$	$j_2$	$k_2$
5	$e_1$	$j_1$	$n_1 + c_2$	$r_1 + h_2$	$t_1 + l_2$	$\bar{u}_1 + m_2$	$n_2$	$o_2$
6	$f_1$	$k_1$	$o_1 + d_2$	$t_1 + l_2$	$\bar{u}_1 + m_2$	$\bar{w}_1 + p_2$	$r_2$	$s_2$
7	0	0	$e_2$	$j_2$	$n_2$	$r_2$	$t_2$	$\bar{u}_2$
8	0	0	$f_2$	$k_2$	$o_2$	$s_2$	$\bar{u}_2$	$\bar{w}_2$

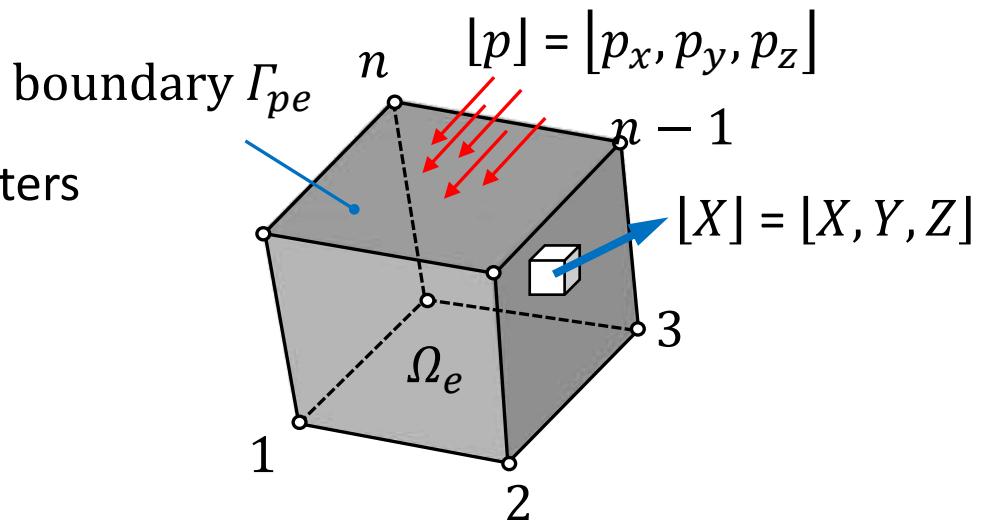
# Potential energy of loading in a finite element

$\{q\}_e$  - local vector of nodal parameters  
 $n_e \times 1$

potential energy of loading  
in a finite element :

$$\begin{aligned}
W_e &= \int_{\Omega_e} [X]\{u\} d\Omega_e + \int_{\Gamma_{pe}} [p]\{u\} d\Gamma_{pe} = \int_{\Omega_e} [X][N]\{q\}_e d\Omega_e + \int_{\Gamma_{pe}} [p][N]\{q\}_e d\Gamma_{pe} = \\
&\quad \underbrace{\{u\} = [N]\{q\}_e}_{\substack{3 \times 1 \\ 3 \times n_e \\ n_e \times 1}} \\
&= (\int_{\Omega_e} [X][N] d\Omega_e + \int_{\Gamma_{pe}} [p][N] d\Gamma_{pe}) \{q\}_e = ([F^X]_e + [F^p]_e) \{q\}_e = [F]_e \{q\}_e
\end{aligned}$$

equivalent load vector:



$$[F]_e = [F^X]_e + [F^p]_e$$

# Equivalent load vector

$$\underset{1 \times n_e}{[F]_e} = \underset{1 \times n_e}{[F^X]_e} + \underset{1 \times n_e}{[F^p]_e}$$

equivalent load vector due to mass forces:

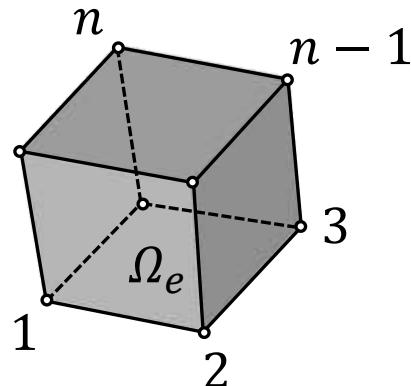
$$\begin{aligned} \underset{1 \times n_e}{[F^X]_e} &= \int_{\Omega_e} [X][N] d\Omega_e = \\ &= \int_{\Omega_e} [X, Y, Z] \begin{bmatrix} N_1 & 0 & 0 & N_2 & 0 & 0 & N_n & 0 & 0 \\ 0 & N_1 & 0 & 0 & N_2 & 0 & \dots & 0 & N_n \\ 0 & 0 & N_1 & 0 & 0 & N_2 & 0 & 0 & N_n \end{bmatrix} d\Omega_e \end{aligned}$$

equivalent load vector due to surface load:

$$\begin{aligned} \underset{1 \times n_e}{[F^p]_e} &= \int_{\Gamma_{pe}} [p][N] d\Gamma_{pe} = \\ &= \int_{\Gamma_{pe}} [p_x, p_y, p_z] \begin{bmatrix} N_1 & 0 & 0 & N_2 & 0 & 0 & N_n & 0 & 0 \\ 0 & N_1 & 0 & 0 & N_2 & 0 & \dots & 0 & N_n \\ 0 & 0 & N_1 & 0 & 0 & N_2 & 0 & 0 & N_n \end{bmatrix} d\Gamma_{pe} \end{aligned}$$

# Potential energy of loading in a finite element

local notation:



$n$  – no. of nodes per FE

$n_p$  – no. of nodal parameters per node

no. of degrees of freedom in FE :

$$n_e = n \cdot n_p$$

$\{q\}_e$  - local vector of nodal parameters  
 $n_e \times 1$

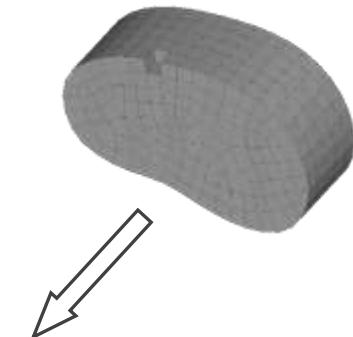
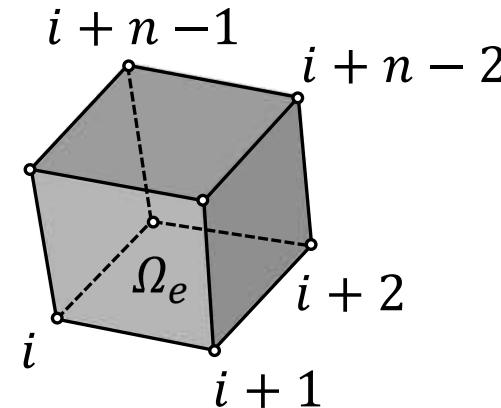
$$W_e = [q]_e \{F\}_e$$

$1 \times n_e \quad n_e \times 1$



equivalent load vector

global notation:



$NON$  – no. of nodes

$n_p$  – no. of nodal parameters per node

no. of degrees of freedom :

$$NDOF = NON \cdot n_p$$

$\{q\}$  - global vector of nodal parameters  
 $NDOF \times 1$

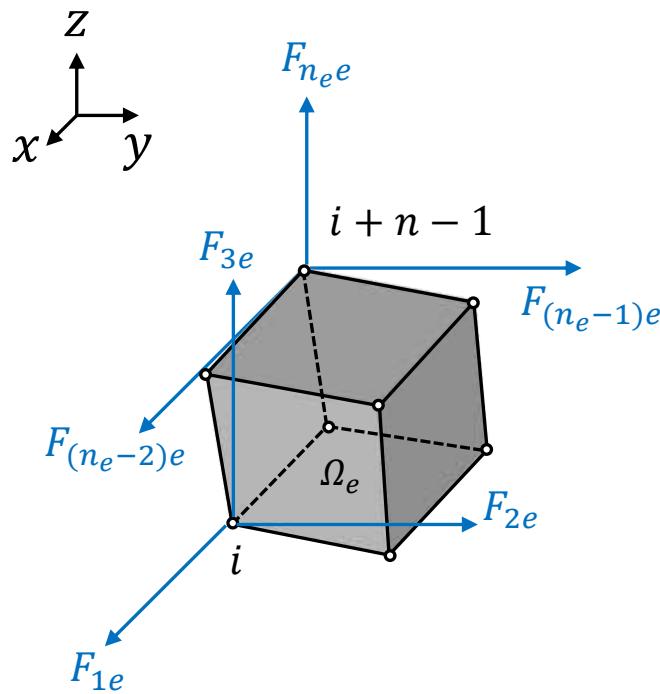
$$W_e = [q] \cdot \{F\}_e^*$$

$1 \times NDOF \quad NDOF \times 1$



extended equivalent load vector

# Extended equivalent load vector in a finite element



extended equivalent load vector:

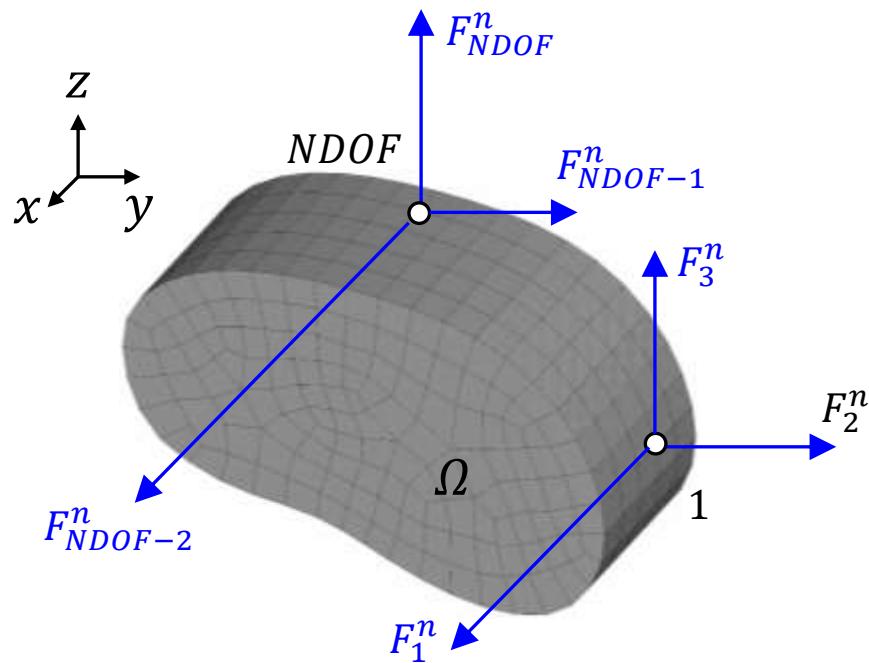
$$\{F\}_e^* = \begin{Bmatrix} 0 \\ 0 \\ \dots \\ 0 \\ F_{1e} \\ F_{2e} \\ \dots \\ F_{n_e e} \\ 0 \\ \dots \\ 0 \end{Bmatrix}_{NDOF \times 1} \quad \begin{array}{ll} 1 & \\ 2 & \\ & \\ j-1 & \\ j & \\ j+1 & \\ & \\ j+n_e-1 & \\ j+n_e & \\ & \\ NDOF & \end{array}$$

equivalent load vector:

$$\{F\}_e = \begin{Bmatrix} F_{1e} \\ F_{2e} \\ F_{3e} \\ \dots \\ F_{(n_e-2)e} \\ F_{(n_e-1)e} \\ F_{n_e e} \end{Bmatrix}_{n_e \times 1}$$

(assumed ascending order  
of components)

# Forces applied directly on nodes. Potential energy of nodal loads



nodal load vector:

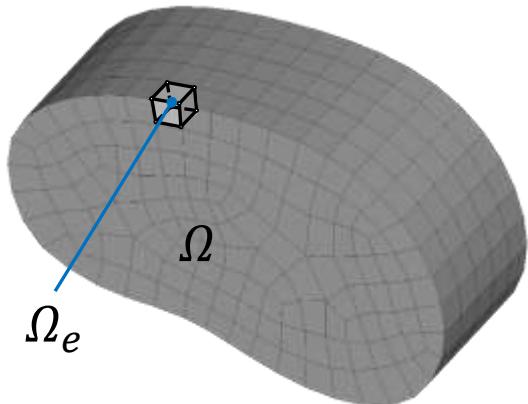
$$\{F\}^n = \begin{Bmatrix} F_1^n \\ F_2^n \\ F_3^n \\ \dots \\ F_{NDOF-2}^n \\ F_{NDOF-1}^n \\ F_{NDOF}^n \end{Bmatrix}_{NDOF \times 1}$$

potential energy of nodal loads:

$$W^n = [q] \cdot \{F\}^n$$

$1 \times NDOF \quad NDOF \times 1$

# Potential energy of loading in a FE model. Global load vector



potential energy of element loads:

$$\Omega = \sum_{e=1}^{NOE} \Omega_e \rightarrow W^e = \sum_{e=1}^{NOE} W_e$$

$NOE$  – no. of FEs

$NDOF$  – no. of degrees of freedom

potential energy  
of nodal loads

$$W = W^e + W^n$$

$$W = \sum_{e=1}^{NOE} W_e + W^n = \sum_{e=1}^{NOE} [q] \cdot \{F\}_e^* + [q] \cdot \{F\}^n = [q] \cdot \left( \sum_{e=1}^{NOE} \{F\}_e^* + \{F\}^n \right)$$

$$= [q] \cdot (\{F\}^e + \{F\}^n) \quad \rightarrow \quad W = [q] \cdot \{F\}$$

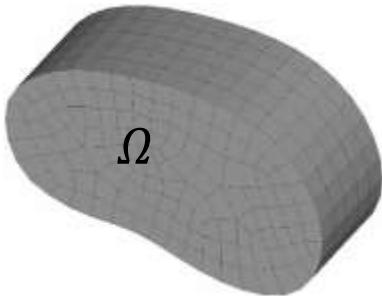
global load vector  
of element loads

nodal load vector

global load vector:

$$\{F\} = \{F\}^e + \{F\}^n$$

# Total potential energy in a FE model. Set of linear equations



Total potential energy of the entire model:

$$V = U - W = \frac{1}{2} \cdot [q] \cdot [K] \cdot \{q\} - [q] \cdot \{F\}$$

$1 \times NDOF$     $NDOF \times NDOF$     $NDOF \times 1$     $1 \times NDOF$     $NDOF \times 1$

$$\{q\} = ?$$

$NDOF \times 1$

$NOE$  – no. of FEs

$NDOF$  – no. of degrees of freedom

$$V \rightarrow \min$$

$$\frac{\partial V}{\partial q_j} = 0 \rightarrow$$

$$[K] \cdot \{q\} = \{F\}$$

$NDOF \times NDOF$     $NDOF \times 1$     $NDOF \times 1$



*set of linear algebraic equations*

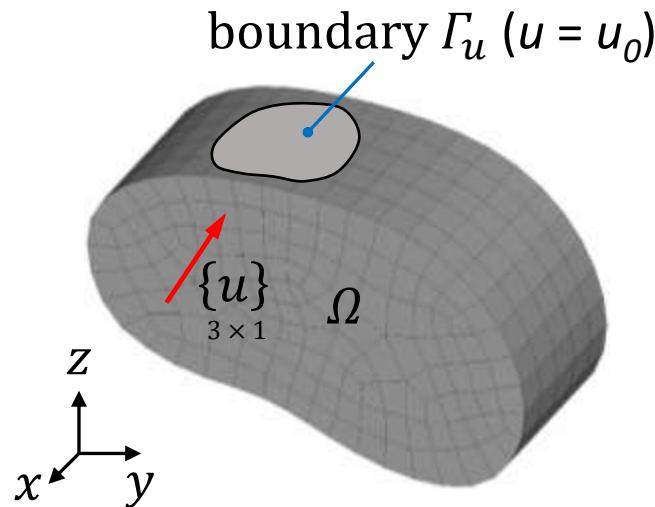
$$\det ([K]) = 0$$

$NDOF \times NDOF$



# Set of FE equations with boundary conditions

The displacement field  $\{u\}$  that represents solution of the problem **fulfils displacement boundary conditons on  $\Gamma_u$**  and minimizes the total potential energy  $V$ .



$NDOF$  – no. of degrees of freedom

$NOF$  – no. of known degrees of freedom on  $\Gamma_u$

$N$  – number of unknown degrees of freedom:

$$N = NDOF - NOF$$

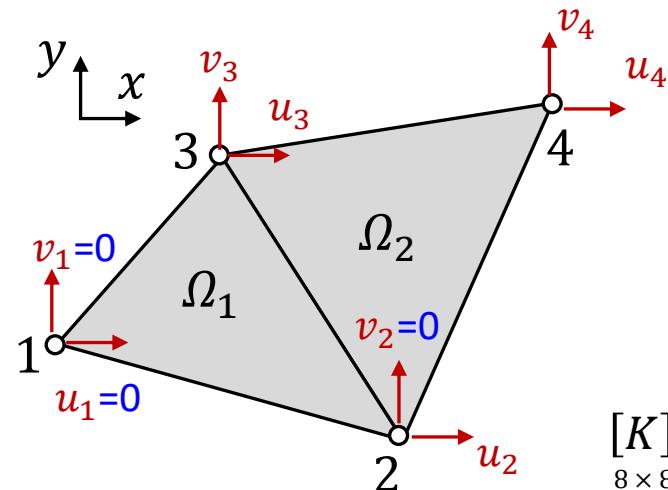
$$\begin{matrix} [K] \\ NDOF \times NDOF \end{matrix} \rightarrow \begin{matrix} [K] \\ N \times N \end{matrix} ; \quad \begin{matrix} \{q\} \\ NDOF \times 1 \end{matrix} \rightarrow \begin{matrix} \{q\} \\ N \times 1 \end{matrix} ; \quad \begin{matrix} \{F\} \\ NDOF \times 1 \end{matrix} \rightarrow \begin{matrix} \{F\} \\ N \times 1 \end{matrix}$$

$$\boxed{\begin{matrix} [K] \\ N \times N \end{matrix} \cdot \begin{matrix} \{q\} \\ N \times 1 \end{matrix} = \begin{matrix} \{F\} \\ N \times 1 \end{matrix}}$$

$$\det \left( \begin{matrix} [K] \\ N \times N \end{matrix} \right) \neq 0$$

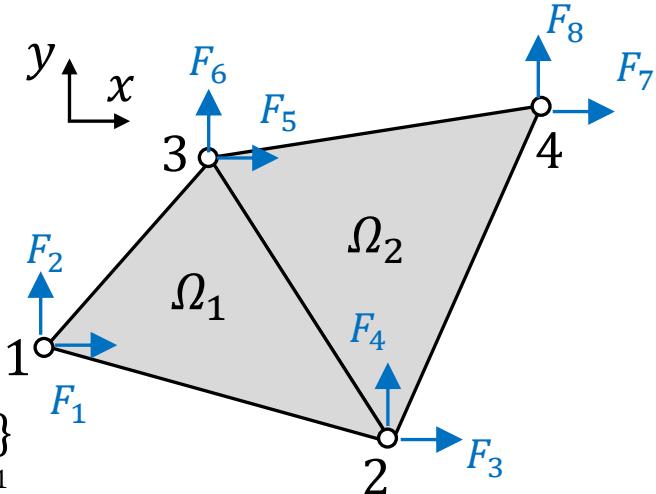
*linear set of algebraic equations with boundary conditions*

## Example 5. Boundary conditions for 2D problem. FE model with two 3-node triangles



$$NDOF = 8 \\ NOF = 3$$

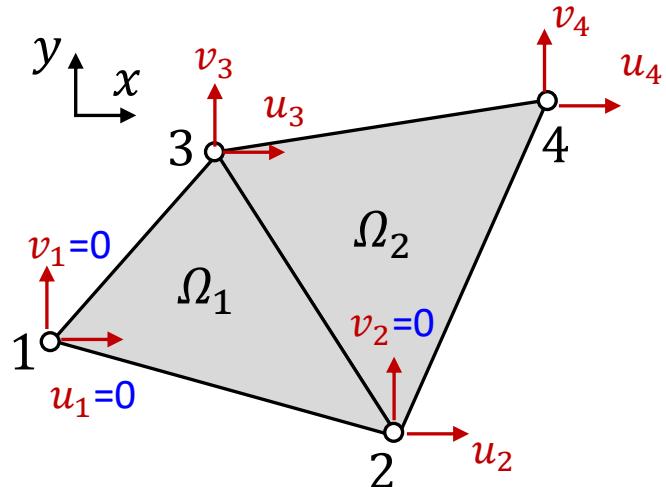
$$[K]_{8 \times 8} \cdot \{q\}_{8 \times 1} = \{F\}_{8 \times 1}$$



1	2	3	4	5	6	7	8
$a_1$	$b_1$	$c_1$	$d_1$	$e_1$	$f_1$	0	0
$b_1$	$g_1$	$h_1$	$i_1$	$j_1$	$k_1$	0	0
$c_1$	$h_1$	$l_1 + a_2$	$m_1 + b_2$	$n_1 + c_2$	$o_1 + d_2$	$e_2$	$f_2$
$d_1$	$i_1$	$m_1 + b_2$	$p_1 + g_2$	$r_1 + h_2$	$s_1 + i_2$	$j_2$	$k_2$
$e_1$	$j_1$	$n_1 + c_2$	$r_1 + h_2$	$t_1 + l_2$	$\bar{u}_1 + m_2$	$n_2$	$o_2$
$f_1$	$k_1$	$o_1 + d_2$	$t_1 + l_2$	$\bar{u}_1 + m_2$	$\bar{w}_1 + p_2$	$r_2$	$s_2$
0	0	$e_2$	$j_2$	$n_2$	$r_2$	$t_2$	$\bar{u}_2$
0	0	$f_2$	$k_2$	$o_2$	$s_2$	$\bar{u}_2$	$\bar{w}_2$

$$\left\{ \begin{array}{l} u_1 = 0 \\ v_1 = 0 \\ u_2 \\ v_2 = 0 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{array} \right\} = \left\{ \begin{array}{l} F_1 \\ F_2 \\ F_3 \\ F_4 \\ F_5 \\ F_6 \\ F_7 \\ F_8 \end{array} \right\}$$

## Example 5. Boundary conditions for 2D problem. FE model with two 3-node triangles

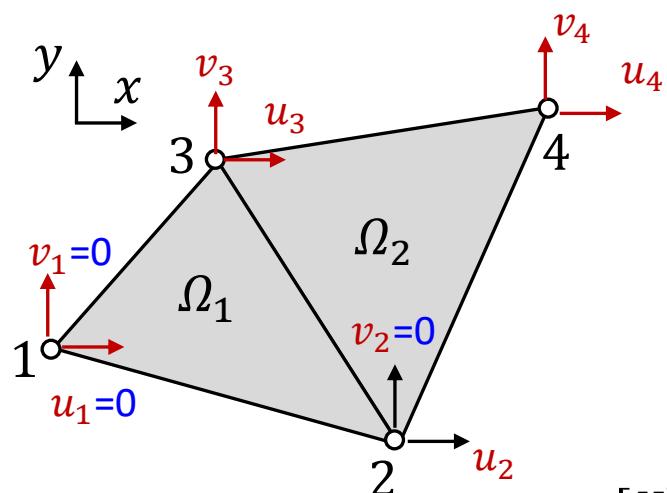


$$[K]_{8 \times 8} \cdot \{q\}_{8 \times 1} = \{F\}_{8 \times 1}$$

1	2	3	4	5	6	7	8	
1	$a_1$	$b_1$	$c_1$	$d_1$	$e_1$	$f_1$	0	0
2	$b_1$	$g_1$	$h_1$	$i_1$	$j_1$	$k_1$	0	0
3	$c_1$	$h_1$	$l_1 + a_2$	$m_1 + b_2$	$n_1 + c_2$	$o_1 + d_2$	$e_2$	$f_2$
4	$d_1$	$i_1$	$m_1 + b_2$	$p_1 + g_2$	$r_1 + h_2$	$s_1 + i_2$	$j_2$	$k_2$
5	$e_1$	$j_1$	$n_1 + c_2$	$r_1 + h_2$	$t_1 + l_2$	$\bar{u}_1 + m_2$	$n_2$	$o_2$
6	$f_1$	$k_1$	$o_1 + d_2$	$t_1 + l_2$	$\bar{u}_1 + m_2$	$\bar{w}_1 + p_2$	$r_2$	$s_2$
7	0	0	$e_2$	$j_2$	$n_2$	$r_2$	$t_2$	$\bar{u}_2$
8	0	0	$f_2$	$k_2$	$o_2$	$s_2$	$\bar{u}_2$	$\bar{w}_2$

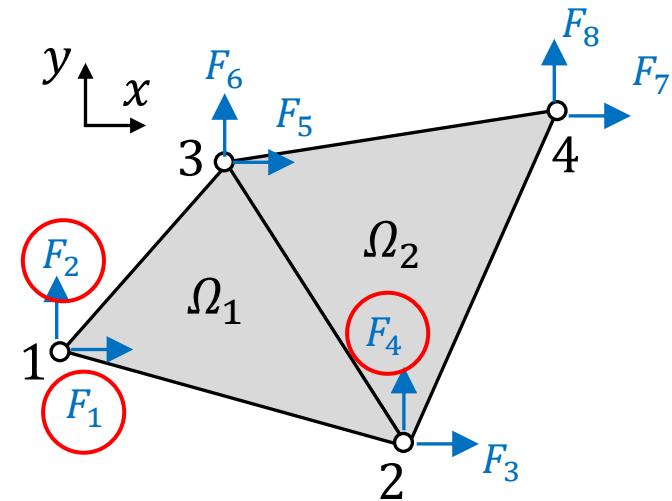
$$\left. \begin{array}{l} u_1 = 0 \\ v_1 = 0 \\ u_2 \\ v_2 = 0 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{array} \right\} = \left. \begin{array}{l} F_1 \\ F_2 \\ F_3 \\ F_4 \\ F_5 \\ F_6 \\ F_7 \\ F_8 \end{array} \right\}$$

## Example 5. Boundary conditions for 2D problem. FE model with two 3-node triangles



$$N = 8 - 3 = 5$$

$$[K]_{5 \times 5} \cdot \{q\}_{5 \times 1} = \{F\}_{5 \times 1}$$



$l_1 + a_2$	$n_1 + c_2$	$o_1 + d_2$	$e_2$	$f_2$
$n_1 + c_2$	$t_1 + l_2$	$u_1 + m_2$	$n_2$	$o_2$
$o_1 + d_2$	$u_1 + m_2$	$w_1 + p_2$	$r_2$	$s_2$
$e_2$	$n_2$	$r_2$	$t_2$	$\bar{u}_2$
$f_2$	$o_2$	$s_2$	$\bar{u}_2$	$\bar{w}_2$

$$\left\{ \begin{array}{c} u_2 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{array} \right\} = \left\{ \begin{array}{c} F_3 \\ F_5 \\ F_6 \\ F_7 \\ F_8 \end{array} \right\}$$

linear set of algebraic equations with boundary conditions

## Solution of a set of FE equations with boundary conditions

$$\begin{matrix} [K] \cdot \{q\} = \{F\} \\ N \times N \quad N \times 1 \end{matrix} \rightarrow \det([K]) \neq 0 \rightarrow \begin{matrix} \{q\} = [K]^{-1} \{F\} \\ N \times 1 \quad N \times N \quad N \times 1 \end{matrix}$$

*DOF solution:*

$$\{q\}$$

*Element solution (ES):*

$$\{\varepsilon\} = [B]\{q\}_e \quad ; \quad \{\sigma\} = [D] \{\varepsilon\} = [D] [B] \{q\}_e$$

$\begin{matrix} 6 \times 1 & 6 \times n_e & n_e \times 1 \\ \uparrow & & \\ \text{a element} & & \end{matrix} \quad \begin{matrix} 6 \times 1 & 6 \times 6 & 6 \times 1 & 6 \times 6 & 6 \times n_e & n_e \times 1 \\ \uparrow & & & & & \\ \text{stress in a finite element} & & & & & \end{matrix}$

*Nodal solution (NS):*

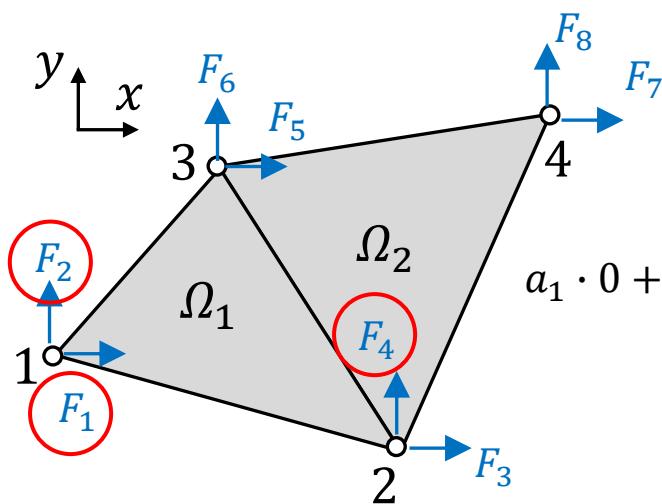
$$(NS)_i = \frac{\sum_{e=1}^k (ES)_{ei}}{k}$$

$(NS)_i$  – avaraged nodal solution at node ( $i$ )

$(ES)_{ei}$  – element solution in element ( $e$ ) and at node ( $i$ )

$k$  – no. of elements adjacent to node ( $i$ )

## Example 6. Reactions calculation for 2D problem. FE model with two 3-node triangles



$$[K]_{8 \times 8} \cdot \{q\}_{8 \times 1} = \{F\}_{8 \times 1}$$

known

$$\boxed{\quad} \cdot \boxed{\quad} = F_1$$

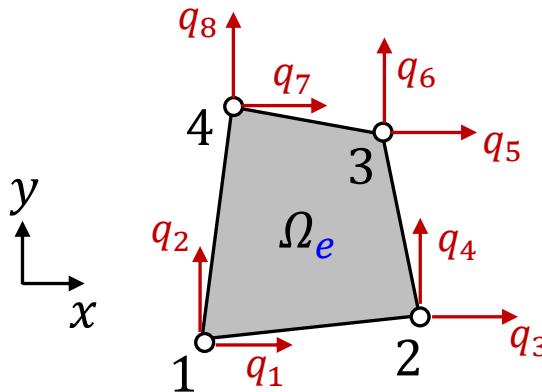
$$a_1 \cdot 0 + b_1 \cdot 0 + c_1 \cdot u_2 + d_1 \cdot 0 + e_1 \cdot u_3 + f_1 \cdot v_3 + 0 \cdot u_4 + 0 \cdot v_4 = F_1$$

$$\boxed{\quad} \cdot \boxed{\quad} = F_2 ; \quad \boxed{\quad} \cdot \boxed{\quad} = F_4$$

	1	2	3	4	5	6	7	8
1	$a_1$	$b_1$	$c_1$	$d_1$	$e_1$	$f_1$	0	0
2	$b_1$	$g_1$	$h_1$	$i_1$	$j_1$	$k_1$	0	0
3	$c_1$	$h_1$	$l_1 + a_2$	$m_1 + b_2$	$n_1 + c_2$	$o_1 + d_2$	$e_2$	$f_2$
4	$d_1$	$i_1$	$m_1 + b_2$	$p_1 + g_2$	$r_1 + h_2$	$s_1 + i_2$	$j_2$	$k_2$
5	$e_1$	$j_1$	$n_1 + c_2$	$r_1 + h_2$	$t_1 + l_2$	$\bar{u}_1 + m_2$	$n_2$	$o_2$
6	$f_1$	$k_1$	$o_1 + d_2$	$t_1 + l_2$	$\bar{u}_1 + m_2$	$\bar{w}_1 + p_2$	$r_2$	$s_2$
7	0	0	$e_2$	$j_2$	$n_2$	$r_2$	$t_2$	$\bar{u}_2$
8	0	0	$f_2$	$k_2$	$o_2$	$s_2$	$\bar{u}_2$	$\bar{w}_2$

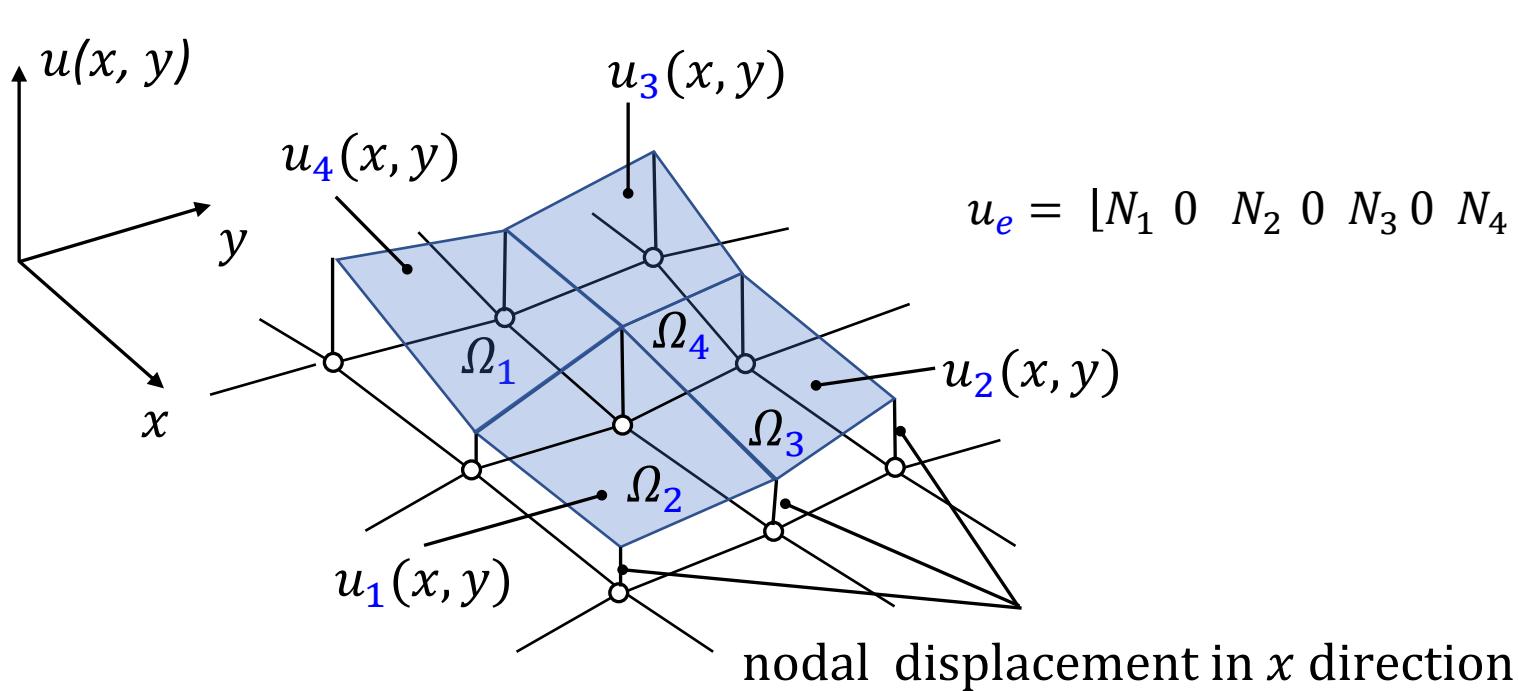
$$\left\{ \begin{array}{l} u_1 = 0 \\ v_1 = 0 \\ u_2 \\ v_2 = 0 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{array} \right\} = \left\{ \begin{array}{l} F_1 \\ F_2 \\ F_3 \\ F_4 \\ F_5 \\ F_6 \\ F_7 \\ F_8 \end{array} \right\}$$

## Example 7. DOF solution $u(x,y)$ for 2D problem. FE model with 4-node quadrilateral elements



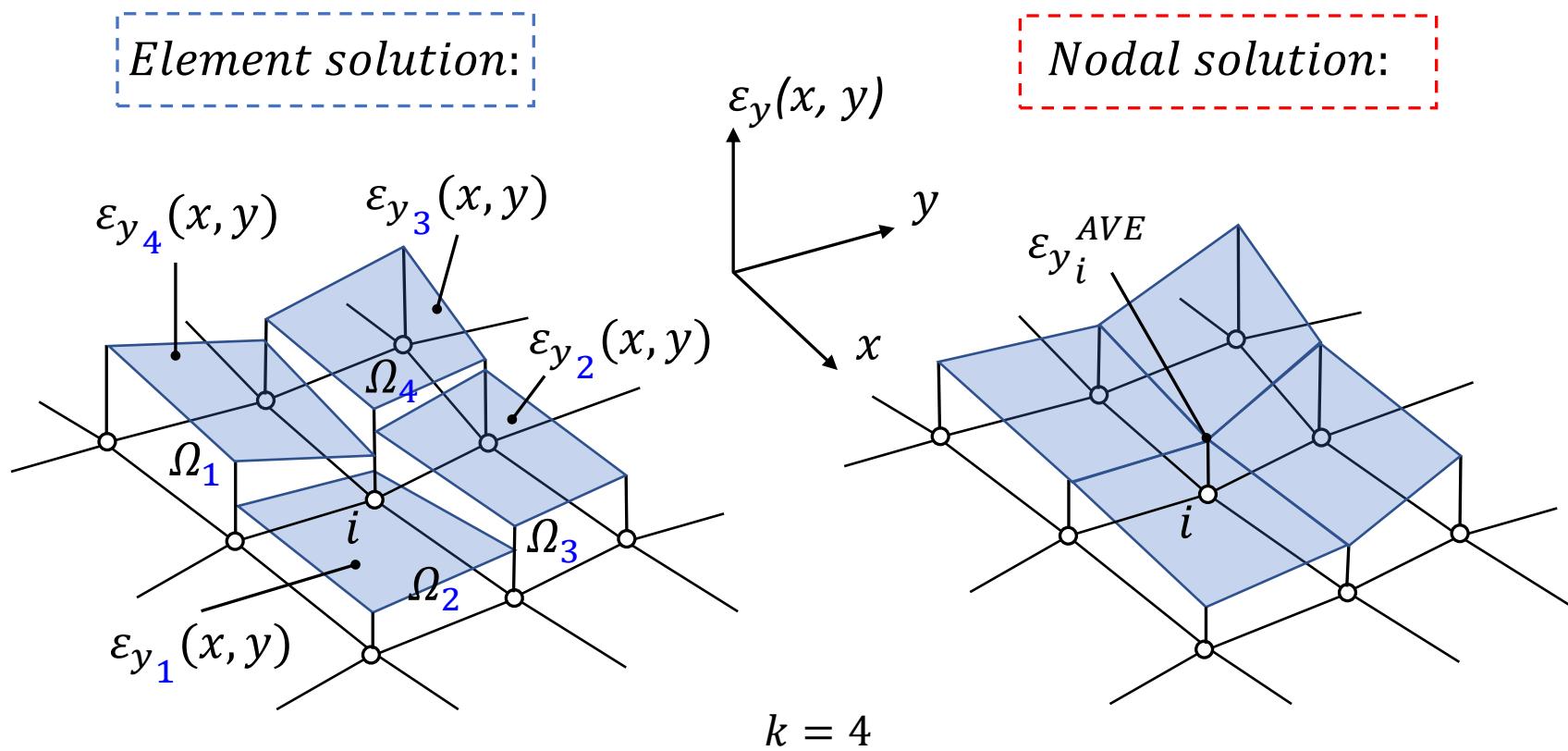
$$\{u\} = [N] \{q\}_e$$

$u_e(x, y)$  – displacement in  $x$  direction



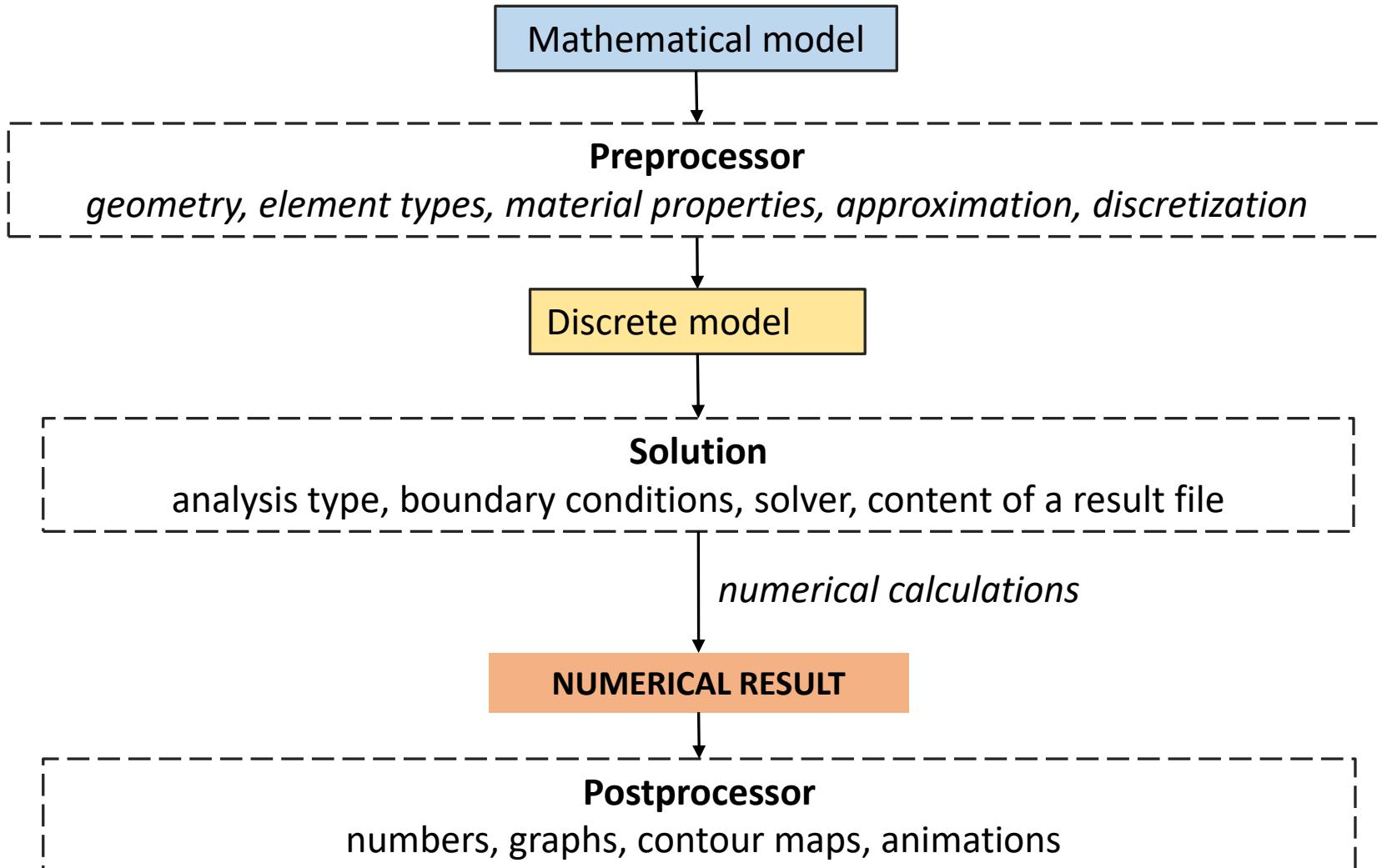
$$u_e = [N_1 \ 0 \ N_2 \ 0 \ N_3 \ 0 \ N_4 \ 0] \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \\ q_5 \\ q_6 \\ q_7 \\ q_8 \end{Bmatrix}_e$$

## Example 8. Strain component $\varepsilon_y(x,y)$ for 2D problem. FE model with 4-node quadrilateral elements

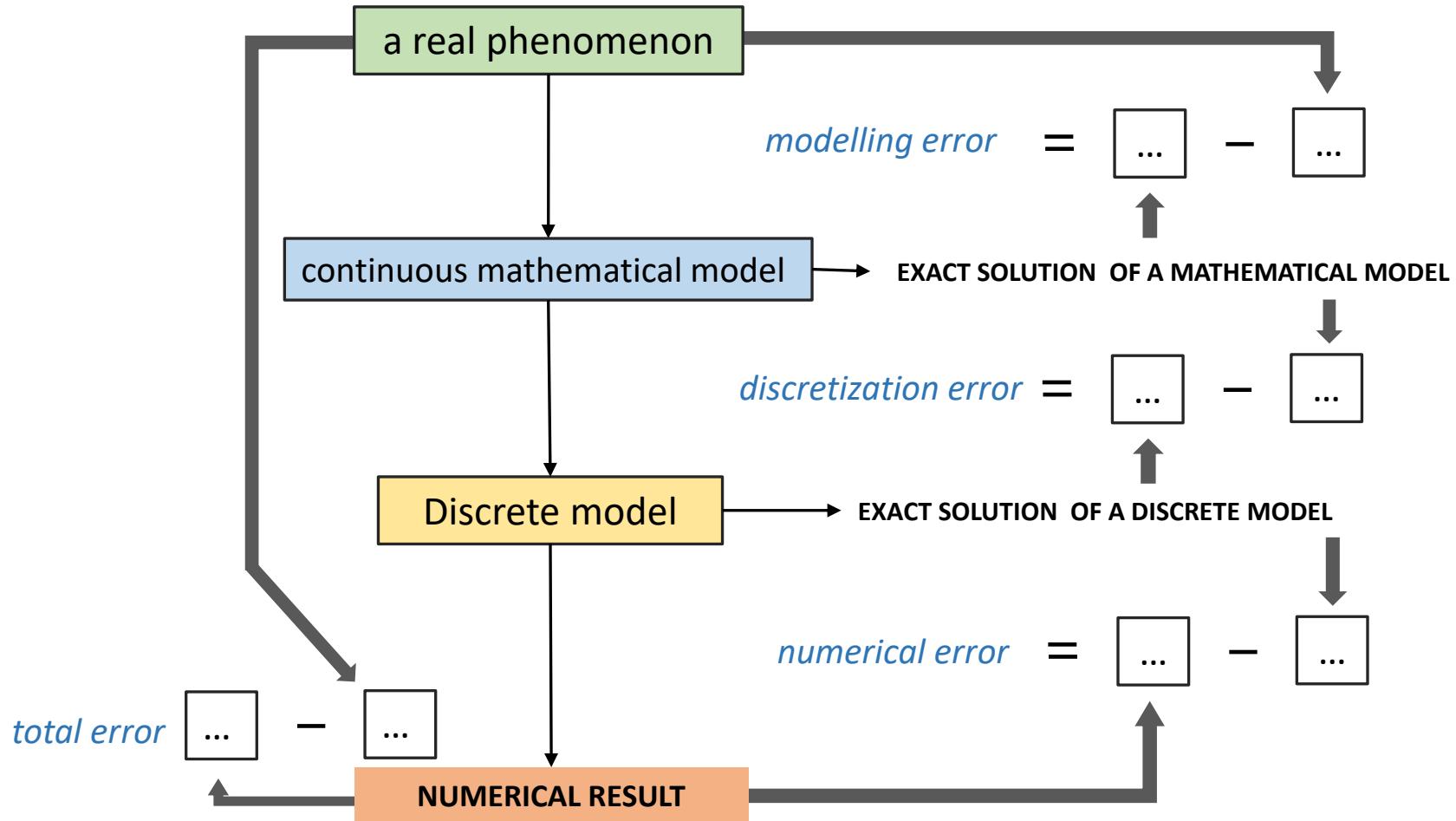


$$\varepsilon_y_i^{AVE} = \frac{\varepsilon_{y_1}(x_i, y_i) + \varepsilon_{y_2}(x_i, y_i) + \varepsilon_{y_3}(x_i, y_i) + \varepsilon_{y_4}(x_i, y_i)}{4}$$

# FE modelling – basic steps



# Accuracy of FEM calculations



$$\text{total error} = \text{modelling error} + \text{discretization error} + \text{numerical error}$$

$$\text{modelling error} \approx \text{discretization error} \approx \text{numerical error} \rightarrow \min$$